



Multilayer co-fired microwave dielectric ceramics in $\text{MgTiO}_3\text{-Li}_2\text{TiO}_3$ system with linear temperature coefficient of resonant frequency

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ABSTRACT

To address the nonlinear variation of resonant frequency in a $\text{MgTiO}_3\text{-CaTiO}_3$ system over a wide temperature range ($-40\text{--}105^\circ\text{C}$), $\text{MgTiO}_3\text{-Li}_2\text{TiO}_3$ system ceramics with multilayer architectures are designed. Trilayer $\text{MgTiO}_3\text{-Li}_2\text{TiO}_3\text{-MgTiO}_3$ ceramics (mass ratio=0.2:0.56:0.2) sintered at 1275°C possess a linear and near-zero temperature coefficient of resonant frequency ($\tau_{f+}=-3.0\text{ ppm}/^\circ\text{C}$ ($25\text{--}105^\circ\text{C}$), $\tau_{f-}=2.7\text{ ppm}/^\circ\text{C}$ ($-40\text{--}25^\circ\text{C}$)). Compared with that of random distribution-type $\text{MgTiO}_3\text{-Li}_2\text{TiO}_3$ ceramics, the $Q \times f$ value of trilayer ceramics is significantly improved ($Q \times f=80,000\text{ GHz}$ ($@7.6\text{ GHz}$)). The electric field distribution is examined by the eigenmode solver of High Frequency Structure Simulator software to analyze the effect of the Li_2TiO_3 content and the stacking scheme on the dielectric properties. The cross-sectional morphologies show a dense intermediate layer formed by ion diffusion, which acts as a glue to bond each layer. This work can be conceived as a new strategy for developing 5G wireless communication components with linear temperature-related resonant frequency ($-40\text{--}105^\circ\text{C}$) and low dielectric loss.

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The global climate is complicated. Winter temperatures can drop to -40°C in some high-latitude regions, while summer temperatures can be higher than 45°C in some African regions, driving a higher demand for temperature-stable communication equipment [1,2]. Microwave dielectric ceramics (MWDCs), the cornerstone of fifth-generation communication technology (5G), are regarded as critical for communication systems to meet this requirement. However, studies on the temperature stability of the resonant frequency of MWDCs have mostly focused on the positive temperature range, whereas the negative temperature range has been largely neglected.

MgTiO_3 ceramic systems, which have outstanding microwave dielectric properties, suitable densities, abundant raw materials and low prices ($\epsilon_r=17$, $Q \times f=160,000\text{ GHz}$ ($@7\text{ GHz}$) and $\rho<3.6\text{ g/cm}^3$), can basically meet the needs of 5G devices, except for the τ_f value that deviates from zero ($\tau_{f-}=-50\text{ ppm}/^\circ\text{C}$) [3–5]. Recently, the two-phase composite method was used in $0.95\text{MgTiO}_3\text{-}0.05\text{CaTiO}_3$ system ceramics, achieving a near-zero τ_f value ($20\text{--}80^\circ\text{C}$) [6,7]. Unfortunately, little attention has been directed toward negative-temperature range stability. In our study, the nonlinear resonant frequency of a $0.93\text{MgTiO}_3\text{-}0.07\text{CaTiO}_3$ ceramic (MT-CT)

varied with temperature, leading to suboptimal temperature stability over a wide temperature range. When the positive τ_f value ($25\text{--}105^\circ\text{C}$) was tuned to zero, the negative τ_f value ($-40\text{--}25^\circ\text{C}$) was approximately $+20\text{ ppm}/^\circ\text{C}$ mainly due to the temperature dependence of the dielectric constant of the CaTiO_3 system.

Recently, Zhou et al. achieved remarkable strides in low-temperature fired Li_2TiO_3 ceramics, and found that the dielectric constant of the Li_2TiO_3 system changes smoothly with temperature, which is essentially different from the CaTiO_3 system [8–10]. To improve the linearity of the resonant frequency of the MgTiO_3 system, Li_2TiO_3 is selected as the additive and better temperature stability is achieved with $\text{MgTiO}_3\text{-Li}_2\text{TiO}_3$ system ceramics. Traditional randomly distributed processes inevitably lead to varied defects due to possible incompatibility between crystal structure, ionic charge and ionic radius. To reduce defects and possible chemical reactions, multilayer ceramics are proposed. Generally, dielectric properties are related to the Li_2TiO_3 content and the stacking scheme. With the aid of a High Frequency Structure Simulator from Ansoft Company (ANSOFT HFSS), we investigated the electric field distribution of different stacking schemes and found that the middle layer played a significant role in overall performance. This multilayer ceramic with linear temperature coefficient of resonant frequency over a wide temperature range and a high $Q \times f$ value can be envisioned as a pioneering material for 5G.

MgTiO_3 , Li_2TiO_3 and CaTiO_3 (MT, LT, and CT, respectively) powders were separately synthesized by the conventional solid-state

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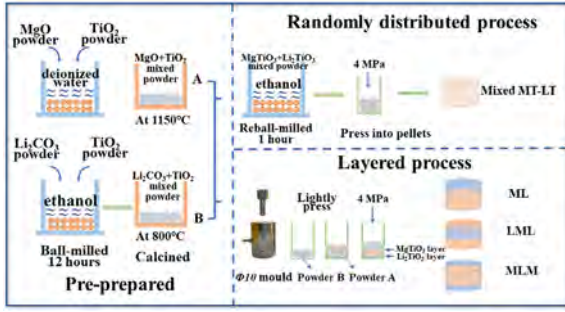


Fig. 1. Schematic diagram of randomly distributed and layered process.

method with high-purity MgO (98%), Li₂CO₃ (98%) and TiO₂ (99%) powders. The powders were weighted according to the desired stoichiometry by a one over ten thousand analytical balance (AW120; Shimadzu Corporation). To achieve an excellent matching characteristic in the diameter direction after co-firing at 1275°C, MgTiO₃ and Li₂TiO₃ powders were calcined at 1150°C and 800°C for 4 h in air, respectively. In addition, CaTiO₃ powders were calcined at 1100°C for 4 h in air.

Fig. 1 shows a graphical representation of the randomly distributed and layered processes for MgTiO₃-Li₂TiO₃ system ceramics. MgTiO₃-Li₂TiO₃-MgTiO₃ (MLM), Li₂TiO₃-MgTiO₃-Li₂TiO₃ (LML) and MgTiO₃-Li₂TiO₃ (ML) ceramics with multilayer architectures were prepared by a layered process. The calcinated MgTiO₃ and Li₂TiO₃ powders were ground with 8 wt% paraffin as a binder. Based on the composition of 0.4MgTiO₃-xLi₂TiO₃ (x=0.5, 0.52, 0.54, 0.56, 0.58), the two powders were alternately added to the mold in three kinds of order. Random distribution-type MgTiO₃-Li₂TiO₃ (mixed MT-LT) was synthesized by a randomly distributed process. The calcinated MgTiO₃ and Li₂TiO₃ powders were mixed according to the mass fraction and then remilled in alcohol for 1 h. In addition, the 0.93MgTiO₃-0.07CaTiO₃ ceramic was also prepared by a randomly distributed process for comparison. The calcinated MgTiO₃ and CaTiO₃ powders were mixed according to the mole ratio and then remilled in deionized water for 1 h. After being dried, the two kinds of mixed powders were ground with 8 wt% paraffin as a binder. All the samples were pressed into pellets with dimensions of 10 mm in diameter and 4–5 mm in thickness under a pressure of 4 MPa for 60 s and then sintered at 1275°C for 4 h in air. The heating rate of the samples was 5 °C/min.

The X-ray diffraction (XRD) patterns of the as-prepared samples were recorded using an X-ray powder diffractometer (D/MAX-2500; Rigaku, Tokyo, Japan) in the range of $10^\circ \leq 2\theta \leq 80^\circ$ using

Cu-K α radiation. Scanning electron microscopy (SEM) and energy dispersive spectrometry (EDS) were used for microstructure observation and cross-sectional analysis (FE-SEM, S-4800; Hitachi, Ltd., Tokyo, Japan). The microwave dielectric performance was measured in TE₀₁₁ mode with a network analyzer (8720ES; Agilent) according to Hakki-Coleman's method and modified by Courtney. The unload Q value was measured by the cavity method. The test system for τ_f values consists of a dielectric resonator (QWED company) and a wide temperature cabinet (GZ-ESPEC GMC-71). The resonant frequencies were obtained by holding the system for 30 min at -40°C, 25°C and 105°C, and the τ_f values were calculated using Eq. (1) and (2):

$$\tau_{f+} = \frac{f_{105} - f_{25}}{f_{25}(T_{105} - T_{25})} \times 10^6 (\text{ppm}/^\circ\text{C}) \quad (1)$$

$$\tau_{f-} = \frac{f_{-40} - f_{25}}{f_{25}(T_{-40} - T_{25})} \times 10^6 (\text{ppm}/^\circ\text{C}) \quad (2)$$

As shown in Fig. 2a, the XRD patterns of the calcined and sintered MgTiO₃ powders show similar peaks that can be indexed to the pure MgTiO₃ phase (JCPDS #06-0494). The detectable reflection peaks of calcined Li₂TiO₃ powders are indexed to the rock-salt monoclinic phase of pure Li₂TiO₃ (JCPDS #33-0831). After the calcined Li₂TiO₃ powders were sintered at 1275°C for 4 h, the intensity of the (002) peak increased while the intensity of the (-133) peak decreased, indicating an increase in the degree of order and better crystallinity [11]. Fig. 2b shows the XRD patterns of the trilayer MLM ceramics with different mass ratios. To confirm the mass ratio of these two phases, the Rietveld refinement method is adopted, and the calculated pattern fits well with the measured pattern as shown in Fig. 2c–g. Fig. 2h shows that the variation trend is similar to the experimental results, but the test values of the Li₂TiO₃ volume fraction are slightly lower, which may be due to the phenomenon of a solid solution between MgTiO₃ and Li₂TiO₃ [12].

Polished and thermal-etched cross-sectional micrographs of the trilayer MgTiO₃-Li₂TiO₃-MgTiO₃ ceramic architecture were investigated by SEM and EDS. Fig. 3a shows a schematic view of the trilayer MLM ceramic, and Fig. 3b shows an SEM image of the corresponding position. The distribution of element and phase composition was analyzed by EDS (see Table. 1). According to the EDS analysis results, spots A and H are Li₂TiO₃ and MgTiO₃. The left side of the SEM image pertains to the Li₂TiO₃ layer, while the right side belongs to the MgTiO₃ layer. Meanwhile, the diffusion of Mg²⁺ and Li⁺ within a certain range generates a dense intermediate layer, which acts as a glue to bond the layers, thus avoiding the aging of

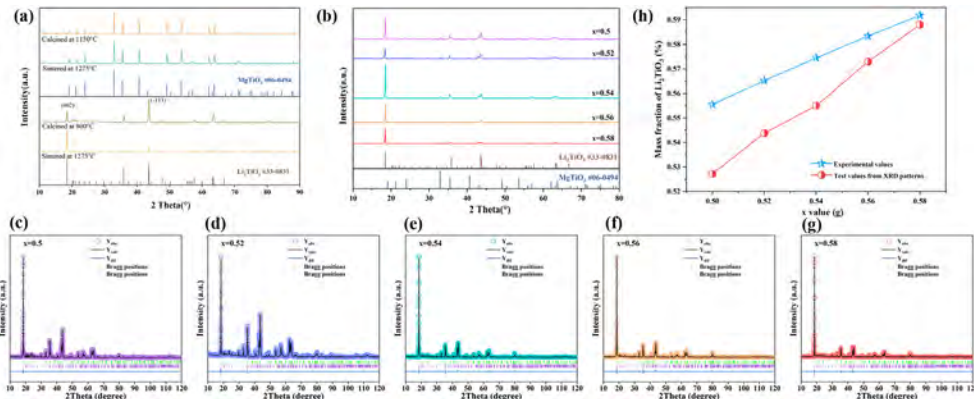


Fig. 2. XRD patterns of as-prepared samples: (a) calcined and sintered MgTiO₃ and Li₂TiO₃ powders; (b) trilayer MgTiO₃-Li₂TiO₃-MgTiO₃ ceramic with different mass ratios; (c)–(g) XRD refinement results; (h) mass fraction of Li₂TiO₃ in the trilayer MgTiO₃-Li₂TiO₃-MgTiO₃ ceramic.

Table 1
EDS data of selected spots marked in the SEM image in Fig. 3.

Spot	Mg (at.%)	Ti (at.%)	O (at.%)	Spot	Mg (at.%)	Ti (at.%)	O (at.%)
A	4.21	25.00	70.78	E	12.97	26.04	61.00
B	2.66	29.63	67.71	F	11.21	31.07	57.72
C	4.29	23.60	72.11	G	12.66	26.16	61.18
D	3.89	25.79	70.32	H	17.64	23.90	58.46

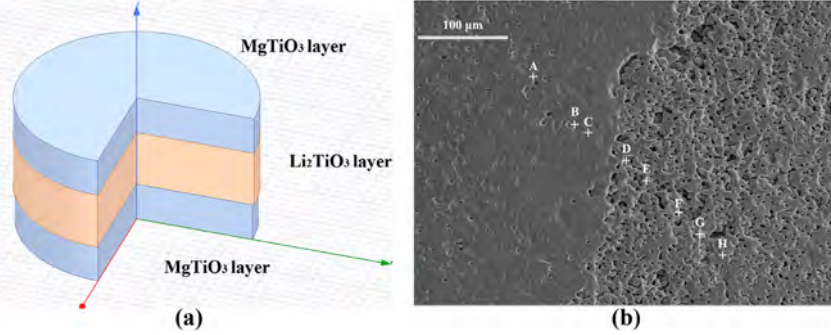


Fig. 3. Lengthwise SEM micrographs for MLM ceramic.

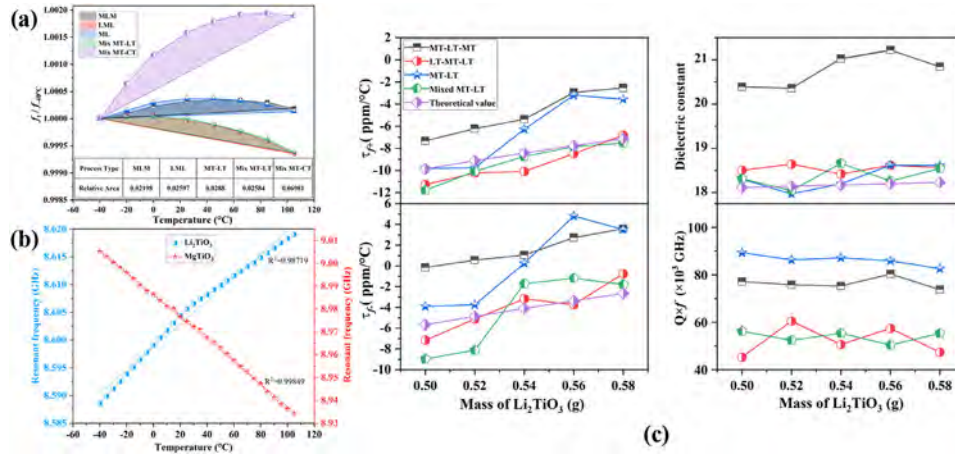


Fig. 4. Variation of (a) R values and (b) resonant frequency with temperature. (c) Microwave dielectric properties of MgTiO_3 - Li_2TiO_3 system ceramics.

traditional organic glue during operation and compensating for the tiny difference in diameter between the two materials.

To illustrate the linear variation of resonant frequency with temperature, the resonant frequencies at multiple temperatures were measured, and the relative resonant frequencies were calculated using the following equation:

$$R = \frac{f_t}{f_{-40^\circ\text{C}}} \quad (3)$$

where f_t is the resonant frequency at each temperature point. Fig. 4a presents the R values of the MgTiO_3 - Li_2TiO_3 system ceramic (mass ratio=0.4:0.56) and the random distribution-type 0.93MgTiO_3 - 0.07CaTiO_3 . The R values of MT-CT exhibit a relatively flat, horizontal curve section in the positive temperature range (25–105°C), indicating a near-zero τ_f value, but a steep curve section in the negative temperature range (−40–25°C). Similar nonlinear variations of f_0 with temperature and the temperature dependence of τ_f are manifested in the Al_2O_3 - TiO_2 system ceramic. Li et al. suggested that these values are determined by the dielectric mixing rules (as shown in Eq. (4)) [13].

$$\tau_{f,C} = \frac{1}{V_1 \varepsilon_{r,1}^\alpha + V_2 \varepsilon_{r,2}^\alpha} [V_1 \varepsilon_{r,1}^\alpha (\tau_{f,1} + \alpha_{L,1}) + V_2 \varepsilon_{r,2}^\alpha (\tau_{f,2} + \alpha_{L,2})] - \alpha_{L,C} \quad (-1 \leq \alpha \leq 1) \quad (4)$$

where the subscripts C, 1 and 2 indicate the composite, Al_2O_3 and TiO_2 phases, respectively, V is the volume fraction, ε_r is the dielectric constant, τ_f is the temperature coefficient of the resonant frequency, α_L is the linear thermal expansion coefficient, and the α value is a determinable value related to the mixing process. From Eq. (4), one can conclude that the temperature-dependent τ_f and ε_r are the main factors affecting the $\tau_{f,C}$ value. Due to the large variation in the dielectric constant of CaTiO_3 with temperature, the τ_f value of the 0.93MgTiO_3 - 0.07CaTiO_3 ceramic is not a constant [14,15]. In turn, the resonant frequency shows nonlinear variation with temperature.

Compared with that of CaTiO_3 , the dielectric constant of Li_2TiO_3 changes more slowly with temperature, and the influence of dielectric constant changes on τ_f is weakened [9,10]. Furthermore, the dielectric constant of Li_2TiO_3 (~19.8) is close to that of MgTiO_3 (~17.0), which is essentially different from the dielectric constant of CaTiO_3 (~170.0). For MgTiO_3 - Li_2TiO_3 system ceramics, $\varepsilon_{\text{MgTiO}_3}$ and $\varepsilon_{\text{Li}_2\text{TiO}_3}$ are approximately equal. Therefore, Eq. (4) can be rewritten as follows:

$$\tau_{f,C} = \frac{1}{V_{\text{MgTiO}_3} + V_{\text{Li}_2\text{TiO}_3}} \left[\frac{V_{\text{MgTiO}_3} (\tau_{f,\text{MgTiO}_3} + \alpha_{L,\text{MgTiO}_3})}{V_{\text{MgTiO}_3} + V_{\text{Li}_2\text{TiO}_3}} + \frac{V_{\text{Li}_2\text{TiO}_3} (\tau_{f,\text{Li}_2\text{TiO}_3} + \alpha_{L,\text{Li}_2\text{TiO}_3})}{V_{\text{MgTiO}_3} + V_{\text{Li}_2\text{TiO}_3}} \right] - \alpha_{L,C} \quad (5)$$

According to Eq. (5), the linear thermal expansion coefficient of the composite ($\alpha_{L,C}$) is assumed to change linearly with the

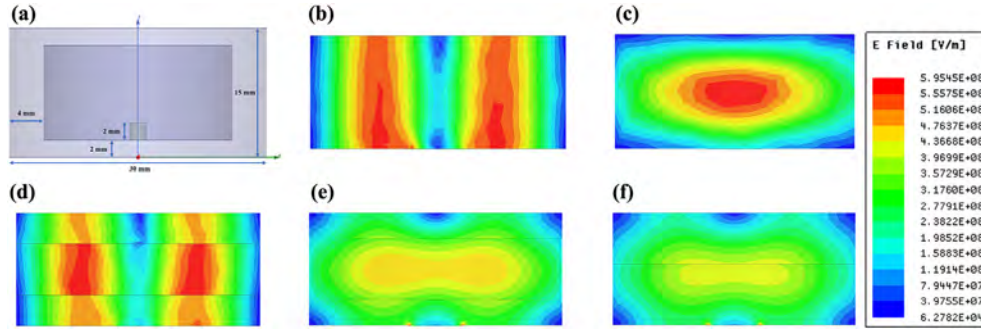


Fig. 5. (a) Diagrammatic drawing of the simulation model and simulation results using the HFSS eigenmode solver in different stacking schemes: (b) MgTiO₃, (c) Li₂TiO₃, (d) LML, (e) MLM, and (f) ML (mass ratio of MgTiO₃ and Li₂TiO₃ = 0.4:0.56).

mass fraction of Li₂TiO₃; thus, the τ_{fC} value is mainly determined by $\tau_{f,MgTiO_3}$ and τ_{f,Li_2TiO_3} . Fig. 4b shows that the resonant frequency changes linearly with the temperature of Li₂TiO₃ and MgTiO₃ (R-squared to the straight line is 0.98719 and 0.99849, respectively), which indicates the relative stability of $\tau_{f,MgTiO_3}$ and τ_{f,Li_2TiO_3} . Thus, the relative frequency curves of the MgTiO₃-Li₂TiO₃ system are practically flattened in the entire temperature range (-40–105°C). Among the several multilayer ceramics with different stacking schemes, the curves are slightly different. We calculated the relative area of the polygon formed by the curve and the baseline (shown in Fig. 4a). Conceivably, the relative area should approach zero when the curve shows satisfactory linearity. The relative areas are in the order of MLM < mixed MT-LT < LML < ML, which indicates better linearity in the MLM ceramics.

The microwave dielectric properties of multilayer ceramics with different stacking schemes and Li₂TiO₃ contents were measured and compared with randomly distributed samples sintered at 1275°C (see Fig. 4b). Simultaneously, the theoretical values of the randomly distributed samples were also given for comparison. The τ_f values of MgTiO₃ and Li₂TiO₃ are -56.3/21.5 ppm/°C in the positive temperature range and -56.6/28.8 ppm/°C in the negative temperature range, respectively. Assuming that MgTiO₃ and Li₂TiO₃ can coexist with each other, the theoretical values (τ_f , ϵ_r) are calculated according to the following empirical formula:

$$\ln \epsilon_r = V_1 \ln \epsilon_{r1} + V_2 \ln \epsilon_{r2} \quad (6)$$

$$\tau_f = V_1 \tau_{f1} + V_2 \tau_{f2} \quad (7)$$

where V_1 and V_2 , ϵ_{r1} and ϵ_{r2} , and τ_{f1} and τ_{f2} are the volume percentages, dielectric constants and temperature coefficients of the resonant frequency of the MgTiO₃ and Li₂TiO₃ phases respectively [10,16]. Compared with the theoretical values, the measured results of mixed MT-LT have similar trends but with slight deviations, which may be due to the difficulty of forming pure two-phase composite ceramics using the conventional mixed route.

In marked contrast, the multilayer ceramics show more positive τ_f values and better $Q \times f$ values. Compared with that of the mixed MT-LT (mass ratio=0.4:0.56), the $Q \times f$ value of trilayer MLM ceramic is significantly improved by approximately 60%. The multilayer architecture could effectively restrict the chemical reaction between MgTiO₃ and Li₂TiO₃ and reduce the influence of defects and second phases on the dielectric properties, resulting in certain differences between the random distribution-type and multilayer ceramics.

By comparison, the dielectric properties of multilayer ceramics with different stacking schemes have similar variation trends. However, different stacking schemes show large variations in τ_f values, $Q \times f$ values and dielectric constants. For example, with 0.56 g Li₂TiO₃, the MLM ceramic shows $\epsilon_r \sim 21.2$, $Q \times f \sim 80,000$ GHz (@7.6 GHz), $\tau_{f+} \sim -3.0$ ppm/°C, and $\tau_{f-} \sim +2.7$ ppm/°C, while the LML ceramic shows $\epsilon_r \sim 18.6$, $Q \times f \sim 57,000$ GHz (@7.7 GHz), $\tau_{f+} \sim -8.5$ ppm/°C, and $\tau_{f-} \sim -3.7$ ppm/°C, and the ML ceramic shows $\epsilon_r \sim 18.6$, $Q \times f \sim 86,000$ GHz (@7.9 GHz), $\tau_{f+} \sim -3.2$ ppm/°C, $\tau_{f-} \sim -4.8$ ppm/°C. Notably, the amount of Li₂TiO₃ required in these multilayer ceramics to achieve a near-zero τ_f is in the order of MLM < ML < mixed MT-LT < LML, indicating that the multilayer architecture plays a certain role in regulating the temperature coefficient and can substantially affect the microwave dielectric properties. To further explore the above differences, the eigenmode solver was used to examine the electric field distributions for different stacking schemes by ANSOFT HFSS.

The finite element method was used to simulate the electric field distribution of the intrinsic resonance of dielectric resonators. Fig. 5a displays a diagrammatic drawing of the simulation model whose size is based on actual values in the experiment. Fig. 5b and Fig. 5c show the electric field distribution of MgTiO₃ and Li₂TiO₃ ceramic resonators, respectively. The heights of MgTiO₃ (3.7727 mm) and Li₂TiO₃ (4.1703 mm) ceramics are fairly close. Excluding the effect of size on simulations, MgTiO₃ and Li₂TiO₃ ceramics showed distinctive electric field distributions. Understandably, the

Table 2
Microwave dielectric properties of typical MgO-TiO₂ and Li₂O-TiO₂ system materials.

Material system	Materials	$Q \times f$ (GHz)	ϵ_r	τ_f (ppm/°C)	Temperature range	Ref.
MgO-TiO ₂ system	(1-x)Mg _(1.04-3/2y) Ce _y TiO ₃ -xCaTiO ₃	71,000	23.0	0.0	20–80°C	[22]
	(1-x)Mg _{0.95} Ni _{0.05} Ti _{0.98} Zr _{0.02} O ₃ -xCa _{0.6} La _{0.8/3} TiO ₃	116,000	23.9	-2.0	20–80°C	[23]
	0.92MgTiO ₃ -0.08NaTaO ₃	55,000	19.4	+1.1	30–90°C	[24]
	(1-x)MgTiO ₃ -x(Ca _{0.6} Na _{0.2} Sm _{0.2})TiO ₃	76,000	22.8	-3.1	25–85°C	[25]
	MgTiO ₃ -5wt%CaF ₂	37,000	17.5	4.8	25–80°C	[26]
Li ₂ TiO ₃ system	(1-x)Li ₂ TiO ₃ -xLi ₂ CeO ₃	59,000	21.2	-7.4	20–80°C	[27]
	0.3Li ₂ TiO ₃ -0.7Li(Zn _{0.5} Ti _{1.5})O ₄	88,000	23.5	-0.8	25–85°C	[28]
	Li ₂ TiO ₃ -20vol.%Li ₂ Zn ₃ Ti ₄ O ₁₂	28,000	22.0	-1.6	25–75°C	[29]
	0.7Li ₂ TiO ₃ -0.3ZnO	100,000	23.0	0.0	25–80°C	[11]
	MLM (mass ratio=0.2:0.56:0.2)	80,000	21.2	-3.0/2.7	-40–105°C	This work

electric field distributions of different stacking schemes were noticeably different. Fig. 5d–f presents the electric field distributions of multilayer ceramics with different stacking schemes (mass ratio of $\text{MgTiO}_3\text{:Li}_2\text{TiO}_3=0.4\text{:}0.56$). In contrast with that of the LML ceramic, the electric field of the ML ceramic tends to concentrate at the bottom due to its asymmetric structure. Li_2TiO_3 shows a positive τ_f value ($\tau_{f+}=21.5$, $\tau_{f-}=28.8$) with a higher ε_r than MgTiO_3 [16–19]. Some studies have indicated that the distribution of the electric field is affected by the distribution of the dielectric constant and is relatively concentrated in the layer with a high dielectric constant, suggesting that the Li_2TiO_3 layer has a denser electric field distribution in the ML structure [16,20,21]. Moreover, since the dielectric constants of MgTiO_3 and Li_2TiO_3 are similar, the electric field distribution of the trilayer architecture is heavily influenced by the middle layer, which plays a more important role in overall performance. Accordingly, one may speculate that the order for the amount of Li_2TiO_3 required to achieve a near-zero τ_f is $\text{MLM} < \text{ML} < \text{LML}$.

The microwave dielectric properties of typical MgO-TiO_2 system and $\text{Li}_2\text{O-TiO}_2$ system materials are summarized in Table 2. The current trilayer MLM ceramic clearly has excellent dielectric properties. In addition, the wide-temperature stability of MgO-TiO_2 system materials was studied and reported for the first time. With the advantage of wide-temperature stability and excellent microwave dielectric properties, the results of this study can be applied to the development of 5G wireless communication components.

In this paper, we selected MgTiO_3 as the matrix, and $\text{MgTiO}_3\text{-Li}_2\text{TiO}_3$ system ceramics with random distribution-type and multilayer architectures were prepared. With a mass ratio of $\text{MgTiO}_3\text{:Li}_2\text{TiO}_3=0.4\text{:}0.56$, the trilayer $\text{MgTiO}_3\text{-Li}_2\text{TiO}_3\text{-MgTiO}_3$ ceramic has excellent properties: $\varepsilon_r=21.2$, $Q \times f=80,000$ GHz (@7.6 GHz) and $\tau_{f+}=-3.0$ ppm/°C, $\tau_{f-}=2.7$ ppm/°C. In contrast to the random distribution-type $\text{MgTiO}_3\text{-Li}_2\text{TiO}_3$, the $Q \times f$ value is significantly improved by approximately 60%. This design provides a stable operation scheme over a wide temperature range (-40–105°C).

Declaration of Competing Interest

None.

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References

- [1] L.Z. Ni, L.X. Li, M.K. Du, J. Alloy. Compd. 844 (2020) 156106.
- [2] L.Z. Ni, L.X. Li, M.K. Du, Y. Zhan, Ceram. Int. 47 (2021) 1034–1039.
- [3] N. Santha, M. Rakhi, G. Subodh, Mater. Chem. Phys. 255 (2020) 123636.
- [4] C.L. Huang, J.L. Hou, C.L. Pan, C.Y. Huang, C.W. Peng, C.H. Wei, Y.H. Huang, J. Alloy. Compd. 450 (2008) 359–363.
- [5] Z.P. Xu, L.X. Li, S.H. Yu, M.K. Du, W.J. Luo, J. Alloy. Compd. 802 (2019) 1–5.
- [6] Y.B. Chen, J. Alloy. Compd. 478 (2009) 657–660.
- [7] H. Li, B. Tang, X. Li, Z.J. Qing, Y.X. Li, H. Yang, Q. Wang, S.R. Zhang, J. Mater. Sci. 49 (2014) 5850–5855.
- [8] H.H. Guo, D. Zhou, L.X. Pang, J.Z. Su, J. Mater. 4 (2018) 368–382.
- [9] H.H. Guo, M.S. Fu, D. Zhou, C. Du, P.J. Wang, L.X. Pang, W.F. Liu, A.S.B. Sombra, J.Z. Su, ACS Appl. Mater. Interfaces 13 (2021) 912–923.
- [10] H.H. Guo, D. Zhou, C. Du, P.J. Wang, W.F. Liu, L.X. Pang, Q.P. Wang, J.Z. Su, C. Singh, S. Trukhanov, J. Mater. Chem. C 8 (2020) 4690–4700.
- [11] C.L. Huang, Y.W. Tseng, J.Y. Chen, J. Eur. Ceram. Soc. 32 (2012) 3287–3295.
- [12] W.J. Luo, L.X. Li, B.W. Zhang, J.L. Qiao, J. Alloy. Compd. 824 (2020) 153978.
- [13] L. Li, S. Yang, S.Y. Wu, X.M. Chen, Appl. Phys. Lett. 118 (21) (2021) 212902.
- [14] H.Y. Zhou, X.Q. Liu, X.L. Zhu, X.M. Chen, J. Am. Ceram. Soc. 101 (5) (2018) 1999–2008.
- [15] V.V. Lemanov, A.V. Sotnikov, E.P. Smirnova, M. Wehnacht, R. Kunze, Solid State Commun. 110 (11) (1999) 611–614.
- [16] J. Zhang, Z.X. Yue, Y. Luo, L.T. Li, Ceram. Int. 44 (2018) 21000–21003.
- [17] H.Z. Zuo, X.L. Tang, H. Guo, Q.G. Wang, C.L. Dai, H.W. Zhang, H. Su, Ceram. Int. 43 (16) (2017) 13913–13917.
- [18] J.J. Bian, Y.M. Ding, Ceram. Int. 40 (2014) 3595–3601.
- [19] Y.Z. Hao, H. Yang, G.H. Chen, Q.L. Zhang, J. Alloy. Compd. 552 (2013) 173–179.
- [20] J. Zhang, Y. Luo, Z.X. Yue, L.T. Li, J. Am. Ceram. Soc. 102 (2019) 342–350.
- [21] S. Wang, W.J. Luo, L.X. Li, M.K. Du, J.L. Qiao, J. Eur. Ceram. Soc. 41 (2021) 418–423.
- [22] W. Gong, B. Ullah, W. Lei, G.F. Fan, X.H. Wang, W.Z. Lu, Ceram. Int. 43 (2017) 3051–3056.
- [23] A. Manan, Z. Ullah, A.S. Ahmad, A. Ullah, D.F. Khan, A. Hussain, M.U. Khan, J. Adv. Ceram. 7 (2018) 72–78.
- [24] K.G. Wang, H.F. Zhou, X.W. Luan, S. Hu, X.J. Zhou, S. He, X. Wang, S.C. Zhou, X.L. Chen, Ceram. Int. 47 (2021) 121–129.
- [25] L.X. Li, S. Li, T. Tian, X.S. Lyu, J. Ye, H. Sun, J. Mater. Sci. Mater. Electron. 27 (2016) 1286–1292.
- [26] J. Zhang, Z.X. Yue, X.H. Zhang, L.T. Li, Ceram. Int. 41 (2015) S515–S519.
- [27] W.Q. Liu, R.Z. Zuo, J. Eur. Ceram. Soc. 38 (2018) 119–123.
- [28] Y. Wu, D. Zhou, J. Guo, L.X. Pang, H. Wang, X. Yao, Mater. Lett. 65 (2011) 2680–2682.
- [29] A. Sayyadi-Shahraki, E. Taheri-Nassaj, S.A. Hassanzadeh-Tabrizi, H. Barzegar-Bafrooei, J. Alloy. Compd. 597 (2014) 161–166.